

Waves in Plasmas

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Recent advances have made great progress towards practical, economical fusion energy

ITER: Access to burning plasma

Advanced plasma control



J-K Park, PPPL



Exascale computing



A. Bhattacharjee

6 × 0 =

Additive/Advanced manufacturing



A. Seltzman, APS-DPP (2020)

Hight temperature

Scientists Just Generated the World's Strongest Magnetic Field

🖒 Like 💭 15 Comments | 🛷 157









We've been making great progress in confining plasmas



S.E. Wurzel and Scott C. Hsu, Physics of Plasmas 29, 062103 (2022)



FIG. 1. Progress towards commercially viable fusion energy requires progress along three equally important axes. This paper focuses only on the axis of energy gain.





Waves References



My <u>SULI talk from a previous year</u>, along with many other cited references in this talk...



Plasmas support wide variety of wave phenomena

- Waves found naturally in plasmas
 - Instabilities, fluctuations, wave-induced transport
- Waves can deliver energy-momentum in plasma
 - Heating, current drive, particle acceleration
 - Mode stabilization, plasma confinement, $\alpha\text{-}$ channeling
- Waves can be used in plasma diagnostics
 - Interferometry, reflectometry, Faraday rotation, Thomson scattering



Photo of aurora: Senior Airman Joshua Strang



First W7-X plasma, IPP, Greifswald 11





Plasmas support wide variety of wave phenomena

• How do we describe waves in plasmas?

• What can the dispersion relation tell us?

• Examples of waves and what we can do with them





- Wave characteristics can change based on surroundings
- Dispersion relation describes relationship between wavelength and frequency of wave, ω(k)











Plasma wave: Similar to pendulum motion





Plasma wave: Similar to ball stuck in a valley





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Plasma wave: Plasma frequency

Compare to a simple harmonic oscillator:



Assume that $\Delta r \sim 1$, then:

$$\frac{1}{\Delta r^2} \sim \Delta r$$







$$m_e \frac{dv}{dt} = -eE$$

$$E \sim \frac{q}{\Delta r^2} = \frac{en_e}{\Delta r^2}$$
$$\frac{d^2r}{dt^2} = -\frac{e^2n_e}{m_e\Delta r^2}$$





- Plasmas respond to magnetic fields, plasmas conduct electricity
- Process to derive plasma dispersion relations
 - 1. Determine assumptions

Is there an external magnetic field and is it uniform?

Are the ions or electrons cold?

What frequency range are you looking at?

Is one species hot?





- Plasmas respond to magnetic fields, plasmas conduct electricity
- Process to derive plasma dispersion relations
 - 1. Determine assumptions
 - 2. Fourier analyze Maxwell's equations to obtain wave equation

Apply Fourier analysis in space and time:

$$\vec{E}, \vec{B} \approx \exp(i\vec{k} \cdot \vec{r} - i\omega t)$$





- Plasmas respond to magnetic fields, plasmas conduct electricity
- Process to derive plasma dispersion relations
 - 1. Determine assumptions
 - 2. Fourier analyze Maxwell's equations to obtain wave equation
 - 3. Obtain dielectric tensor, relates plasma current to electric field
 - 4. Combine the above to yield the dispersion relation, $\omega(k)$

Dielectric tensor – *contains all of the plasma physics*

 $\vec{n} \times \vec{n} \times \vec{E} + \vec{\bar{K}} E = 0$

For **K** = 1, get vacuum waves



Index of refraction: $\bar{n} = \frac{ck}{ck}$



- Plasmas respond to magnetic fields, plasmas conduct electricity
- Process to derive plasma dispersion relations
 - 1. Determine assumptions
 - 2. Fourier analyze Maxwell's equations to obtain wave equation
 - 3. Obtain dielectric tensor, relates plasma current to electric field
 - 4. Combine the above to yield the dispersion relation, $\omega(k)$
- This process only results in waves in plasmas
 - No resulting instabilities because there are no sources

()

- Provides basic framework for how more complex dispersion relations are derived

$$\vec{n} \times \vec{n} \times \vec{E} + \overline{\vec{K}} \bullet E = 0$$

Index of refraction: $\vec{n} = \vec{n}$

For **K** = 1, get vacuum waves





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Dispersion relation contains lots of information

- Wave characteristics can change based on surroundings
- Dispersion relation describes relationship between wavelength and frequency of wave, $\omega(k)$
- Resonances can occur when $n \to \infty$
 - At this point, the wavelength goes to zero
 - The waves resonant with particle motion, phase velocity goes to zero and energy can be absorbed
- Cutoffs occur when $n \rightarrow 0$
 - At this point, the wavelength goes to infinity
 - The waves will not propagate and typically significant reflection occurs







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Interactions of waves and particles

- Collisionless damping allows energy exchange between plasma and electromagnetic waves
- Particles with speed comparable to v_{ph} speed can resonate
- Particles with speed slightly slower than v_{ph} will be accelerated, take energy from wave
- Particles with speed slightly faster than v_{ph} will decelerate, give energy to wave





Waves in cold plasma dispersion relation

- Several waves can be solved for with the cold plasma dispersion relation
 - Only a few examples will be shown in this talk
- Propagation parallel to B₀
 - Plasma oscillations (first example)
 - Alfven waves
 - Whistler waves
- Propagation perpendicular to B₀
 Ordinary and extraordinary waves





Low frequency MHD waves – Alfvén waves

- Very low frequency waves ($\omega \ll \Omega_{Cl}$)
- Magnetohydrodynamic (MHD) wave where ions oscillate in response to a restoring force provided by an effective tension on the magnetic field lines
 - Linearize MHD equations to obtain shear Alfvén
 - EM waves that propagate along magnetic field lines





Alfvén eigenmodes (AE) can cause fast-ion transport

- Fast ions created through injecting energetic neutral particles, ion cyclotron resonance heating, or fusion reactions
- AEs are MHD instabilities driven by wave particle interactions
- In DIII-D, high beam power can drive strong AE activity, causing fast-ion profile to flatten

[Heidbrink et al., PRL 99, 245002 (2007)]









Fast-ion transport can reduce fusion performance and lead to losses that damage fusion reactor walls

- AEs cause transport that can:
 - Reduce absorbed beam heating power
 - Reduce current drive
 - − Reduce achievable $β_N$ (fusion power ∝ ($β_N$)²)
 - Cause fast ion losses that damage walls
- A 'sea' of AEs are predicted to be unstable in ITER
- Important questions:
 - When is transport significant?
 - What can we do to control AE transport



C. Collins AAPS (2018)





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Whistler waves found in magnetosphere

- Originally observed by radio/telephone operators in WWI/II
- Lightning strikes excite broad range of radio frequency waves in magnetosphere
- Some whistlers born at strike site, propagate along earth's dipole field
- Because of dispersion, higher frequency waves go faster than lower frequency
 - Phase and group velocity depend on frequency



THE UNIVERSITY

Whistler waves observed in tokamaks

- Runaway electrons provide driving energy for whistler waves
 - Increasing B suppresses whistlers
 - Decreasing B enhances whistlers
- Observed more whistlers with increased intensity in measured hard x-rays
 - Dispersion relationship suggests electron energy ~10-15 MeV





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Electron cyclotron range waves provide heating, drive current

- Electron cyclotron (EC) frequency is the natural frequency of rotation of electrons in magnetic fields $f_{ce} = \frac{eB}{2\pi m_e c}$
- Considering this frequency range, there are two solutions to the cold plasma dispersion relation
- Ordinary mode (O-mode)
 - E is parallel to B
 - Independent of B
 - Depends on density
- Extraordinary mode (X-mode)
 - E is perpendicular to B
 - Depends on B, n_e







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Both O-mode and X-mode have resonances at the electron cyclotron frequency, cutoffs depend on plasma density





EC waves provide localized heating/current drive

- Launched radiofrequency (RF) waves absorbed near cyclotron resonance
 - Tune to either electron or ion cyclotron motion
 - RF source frequency can be chosen to heat precise radius
 - For tokamaks, $B_t \propto 1/R$







EC waves provide localized heating/current drive





EC waves provide localized heating/current drive

- Many examples of ECH/ECCD in tokamaks and other confinement devices
 - Large-scale, high-performance devices depend on waves for heating
- EC heating/ EC current drive can provide current profile tailoring in TCV
 - Improve central electron energy confinement
 - Stabilize MHD modes

Z.A. Pietrzyk PRL 86, 8 (2001)





Electron cyclotron wave injection provides plasma heating, current drive – in certain conditions

- If plasma is too dense, O-mode and X-mode reflected near plasma edge
 - Happens in spherical tokamaks and stellarators

 $\omega_{\text{source}} > \omega_{pe}$

• Alternative heating method required





Electron Bernstein waves can travel in high density plasmas

- Electron Bernstein Waves (EBW) can only travel inside the plasma
 - Wave moves due to coherent motion of charged particles
- Can only couple to EBW by launching O- or Xmodes







Electron Bernstein waves can propagate in overdense plasmas

- Electron Bernstein waves (EBW) are hot plasma waves:
 - Longitudinal, electrostatic waves
 - Propagates perpendicular to B
 - Do not experience a density cutoff in the plasma



 Cannot propagate in vacuum -> must launch O- or X-mode to mode couple to EBW

$$1 - 2\sum_{s} \frac{4\pi n_{s} m_{s} c^{2}}{\lambda B_{0}^{2}} \left| \sum_{s} e^{-\lambda} I_{n}(\lambda) \frac{n^{2}}{\left(\omega_{\Omega}\right)^{2} - n^{2}} \right| = 0 \quad \text{Where: } \lambda = \frac{k_{\perp}^{2} \kappa T_{\perp}}{m \Omega^{2}}$$

• As wave frequency approaches EC harmonic, $\omega = n\Omega_{C}$, wave is strongly absorbed 62 <u>Wisconsin</u>



EBW emission can be used to measure temperature

- Electron Bernstein wave emission at blackbody levels, proportional to local ${\rm T_e}$





Coupling microwave power to high density fusion plasmas can be difficult - but possible

- Plasma naturally emits microwaves from cyclotron resonance location
- Assumed physics of microwave emission from high density plasmas same as launching
 - Measurements on NSTX didn't agree with predictions
 - Plasma edge had too many collisions, absorbed microwaves
- Unexpected results present opportunities





PEGASUS-III – lighting a match for fusion

- Future spherical tokamaks call for solenoid-free operation
 - Need to minimize solenoid due to shielding/cost
- Solenoid removal simplifies tokamak design
 - Potential cost reduction
 - More space for inboard shielding/blanket
 - Lower electromechanical stresses



PEGASUS Website: Publications, Presentations http://pegasus.ep.wisc.edu





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- Major upgrade to Pegasus facility underway to compare/contrast/combine startup techniques
 - Local helicity injection
 - Coaxial helicity injection
 - EBW startup and sustainment
- University-class fusion facilities provide innovative approaches to fusion energy development

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- Waves can drive turbulence...



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Contact me, Steffi Diem (<u>sjdiem@wisc.edu</u>), for research and educational opportunities on PEGASUS-III

